

Ultra-thin stack of n-type hydrogenated microcrystalline silicon and silicon oxide front contact layer for rear-emitter silicon heterojunction solar cells

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ABSTRACT

We demonstrate the clear advantage of a n-type hydrogenated microcrystalline silicon (n- $\mu\text{c-Si:H}$) seed layer on the optoelectronic properties and crystallisation behaviour of n-type hydrogenated microcrystalline silicon oxide (n- $\mu\text{c-SiO}_x\text{:H}$) front contact layers. The presence of a non-oxidic n- $\mu\text{c-Si:H}$ seed layer can reduce the thickness and refractive index of the n- $\mu\text{c-SiO}_x\text{:H}$ front layer significantly while maintaining a high degree of crystallisation and excellent conductivity. This leads to increase in short-circuit current density (J_{sc}) by 2.64% and open-circuit voltage (V_{oc}) by 0.56% in comparison to that of a device without the seed layer. The enhancement in J_{sc} can be attributed to the reduction in parasitic absorption loss in the extremely thin front layer. In addition, the improvement in V_{oc} can result from enhanced surface passivation of the wafer due to seed layer growth in very high hydrogen plasma environment which can play a role as the hydrogen post-plasma treatment. The low thickness of the n- $\mu\text{c-SiO}_x\text{:H}$ front layer yields lower internal recombination losses. In conjunction with an optimised n- $\mu\text{c-Si:H}$ seed layer and n- $\mu\text{c-SiO}_x\text{:H}$ front layer, we obtained a high conversion efficiency value of 21.8% with V_{oc} of 727 mV, J_{sc} of 39 mA/cm², and FF of 77% among the fabricated cells in laboratory.

1. Introduction

Silicon heterojunction (SHJ) solar cells owe their progress in performance to excellent passivation effects on both sides of the crystalline silicon wafer surface. SHJ solar cells have been in the research forefront for more than twenty years due to the following advantages: a) simple and cost-effective production process without the need for patterning steps, b) high open-circuit voltage (V_{oc}) in comparison with the traditional homojunction crystalline silicon cells, and c) good performance stability [1,2]. High efficiencies far over 24% have been reported for the SHJ cell until now [3–5]. Numerous efforts have been made to improve the SHJ solar cells further by focusing on the following aspects: 1) reducing carrier recombination loss by high-quality surface passivation layer [6,7]; 2) reducing parasitic absorption loss at the front side by optimising the front grid silver electrode with higher aspect ratio and conductivity and/or developing ultra-thin wide-gap front contact layers; and, 3) reducing the material cost and carrier collection loss by using thinner wafers [8,9]. Among these, losses due to optical parasitic absorption at the front contact layer pose an enormous challenge which

limits the performance of SHJ solar cells [10,11].

Generally, in the front junction configuration, p-type doped layers are situated on the front side to collect those minority charge carriers that are generated close to their collecting contact in an efficient manner. This is really necessary in the case of short-lifetime absorber materials with small carrier diffusion lengths of minority charge carriers [1,12]. However, with good surface passivation of hydrogenated amorphous silicon (a-Si:H) materials as also high-quality absorber wafers available today, the use of a p-type front contact layer in SHJ solar cells is less demanding as the minority charge carriers can diffuse effectively to their collecting electrode regardless of the location of the electrode [13]. In addition, a Schottky barrier is formed between the p-type front contact and transparent conducting oxide (TCO) layers due to the inherent weak conductivity of the front surface field layer [14,15]. A thick and/or highly doped p-layer, formed at the expense of optical transparency, is normally able to eliminate this barrier and thus forming a suitable front surface field for hole collection. However, this results in parasitic absorption and carrier recombination at the junction which adversely affects the solar cell parameters, V_{oc} , fill factor (FF),

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and short-circuit current density (J_{sc}). It appears that the contact formation for hole collection at the front side is a delicate task [16]. Efforts made in this direction saw the evolution of a new SHJ device architecture, in which the p-contact layer is located on the rear side, popularly known as the rear-emitter silicon heterojunction (RE-SHJ) solar cell [17,18].

In RE-SHJ configuration, a thicker and/or highly conducting p-contact layer can be used at rear side regardless of any major transparency concerns. The use of the RE-SHJ design relaxes the strict requirement of the front TCO material and, at the same time, avails more optical transmission into the bulk absorber (wafer) through a very thin, n-type front layer without seriously impairing conductivity [13,19]. Standard RE-SHJ solar cells use n- and p-doped hydrogenated amorphous silicon (n-a-Si:H and p-a-Si:H) as carrier-selective front and back contact layers, respectively [20]. The efficient internal electric field which helps to separate and collect carrier charges depends strongly on the net doping of these doped layers. However, a highly uniform degree of doping is hardly ever achieved with a-Si:H material because large dopant quantities cause accumulation of defect density. This has a detrimental effect on uniform doping and even likely to deteriorate the passivation quality at wafer interfaces [21]. Additionally, the low bandgap of the a-Si:H material hinders the transparency of the solar spectrum in the short wavelength region. A material with wide bandgap and high conductivity appears most suitable to alleviate this problem.

Hydrogenated microcrystalline silicon ($\mu\text{-Si:H}$) is found to be highly attractive as a front contact layer in thin film silicon solar cells [22,23] as well as SHJ solar cells [24,25]; this is owing to its high doping efficiency thanks to a crystalline phase and low absorption coefficient in the short wavelength region in comparison to the a-Si:H counterpart. By incorporating oxygen, a bi-phase microcrystalline silicon oxide ($\mu\text{-SiO}_x\text{:H}$) material with columnar crystalline silicon phase mixed in an amorphous silicon oxide matrix phase and considerably improved transparency can be obtained without significantly affecting vertical conductivity [13,24,26]. In such a bi-phase material, high vertical conductivity is achieved throughout crystalline phase while the amorphous oxide phase is responsible for enhanced transparency. A challenge in realising $\mu\text{-SiO}_x\text{:H}$ as the front contact layer in SHJ solar cells is the formation of ultra-thin layers with high-crystalline quality and efficient doping without impairing the a-Si:H passivation layers below. It requires delicate control of the initial nucleation processes to limit the incubation phase during crystal growth of n- $\mu\text{-SiO}_x\text{:H}$ layer.

In this work, we investigate a stack of front contact layers for RE-SHJ solar cells including a very thin n- $\mu\text{-Si:H}$ layer grown under high hydrogen dilution and at low plasma power density in a favourable nucleation environment for subsequent growth of n- $\mu\text{-SiO}_x\text{:H}$ layers. This study helped us to evaluate the optimum thickness of n- $\mu\text{-SiO}_x\text{:H}$ layers that are endowed with high crystalline fraction and electrical conductivity. The analysis of solar cell parameters, namely short-circuit current density (J_{sc}) in the short wavelength region, open-circuit voltage (V_{oc}), and fill factor (FF) based on a stack of seed-crystalline layer and ultra-thin n- $\mu\text{-SiO}_x\text{:H}$ front layer is presented.

2. Experiment

Single layers of n- $\mu\text{-SiO}_x\text{:H}$ with various thicknesses were deposited on EAGLE XG™ glass substrates using a cluster multi-chamber plasma enhanced chemical vapor deposition (PECVD) system. To fast prompt nucleation for n- $\mu\text{-SiO}_x\text{:H}$ growth, the substrate was firstly deposited with n- $\mu\text{-Si:H}$ (seed layer) of 3 nm thickness in the gas mixture consisting of silane (SiH_4), hydrogen (H_2), and phosphine (PH_3). After the n- $\mu\text{-Si:H}$ deposition, the plasma discharge power was interrupted. Then, the second gas mixture including of SiH_4 , H_2 , PH_3 and carbon dioxide (CO_2) was set up for subsequent n- $\mu\text{-SiO}_x\text{:H}$ growth. The deposition parameters of these layers are summarized in Table 1. Optical properties including absorption coefficient (α), and reflective index (n) of n- $\mu\text{-SiO}_x\text{:H}$ layers with and without the seed layer were extracted by

Table 1

Deposition parameters of n- $\mu\text{-Si:H}$ seed layer and n- $\mu\text{-SiO}_x\text{:H}$ front layers.

Deposition Condition	n- $\mu\text{-Si:H}$	n- $\mu\text{-SiO}_x\text{:H}$
$[\text{H}_2]/[\text{SiH}_4]$ gas flux ratio	160	125
$[\text{PH}_3]/[\text{SiH}_4]$ gas flux ratio	0.6	2
$[\text{CO}_2]/[\text{SiH}_4]$ gas flux ratio	0	0.5
Power density (mW/cm^2)	42	42
Working Pressure (Torr)	1.5	1.5
Deposition temperature ($^\circ\text{C}$)	110	110
Thickness (nm)	3	4–17

fitting spectroscopic ellipsometry (SE - VASE® Ellipsometer - J. A. Woollam Co., Inc.) data using Tauc-Lorentz model in the wavelength range from 240 to 1700 nm at incident angle of 65° . The lateral conductivity of these samples was measured using evaporated aluminium coplanar bi-electrodes (250- μm spacing) with a programmable Keithley 617 electrometer. Structural characteristics were evaluated by Raman spectroscopy (Ramboss 500i, Dongwoo Optron Co., Ltd.).

The RE-SHJ solar cells were fabricated on random pyramid structure obtained on a textured n-type Czochralski wafer of 140 μm thickness and 4.3 Ωcm resistivity as shown in the schematic diagram in Fig. 1. The cell area was calibrated by indium-doped tin oxide (ITO) membrane deposited using a mask area of 10.24 cm^2 . After screen printing a Ag-grid electrode on the front side and Ag back electrode using commercial paste (KYOTO) at 160 $^\circ\text{C}$, the fabricated solar cells were characterised by I–V curves in dark and illuminated conditions measured under standard test conditions (one-sun illumination, 25 $^\circ\text{C}$, AM1.5 spectrum, 100 mW/cm^2 incident light power) through 10.24 cm^2 shadow mask. The width and height of the front Ag-grid electrode were optimised as 40 μm and 20 μm , respectively, to reduce the shadowing down to $\sim 4\%$. In addition, the external quantum efficiency (EQE) was measured using a solar cell quantum efficiency measurement system (QEX7, PV Measurement Inc.) in the wavelength range of 300–1100 nm.

3. Result and discussion

First, we estimated the optoelectronic and structural characteristics of single n- $\mu\text{-SiO}_x\text{:H}$ layers of various thicknesses deposited in the following geometry: glass substrate/seed layer of 3 nm/n- $\mu\text{-SiO}_x\text{:H}$ layer (varies from 5 to 17 nm). The characteristics of single n- $\mu\text{-SiO}_x\text{:H}$ layers without the seed layer were derived for comparison. The optoelectronic properties of these samples, such as the lateral conductivity (σ) and refractive index (n) at 400 and 632 nm are shown in Table 2. Fig. 2 shows plots of $\{\alpha(E)E\}^{1/2}$ versus photon energy E as used for the determination of optical band-gap of the films. It shows that all n- $\mu\text{-SiO}_x\text{:H}$ films have a wide optical band gap around 2.1–2.4 eV. The structural analyses of such layers using Raman spectroscopy are shown in Fig. 3a and b. Crystalline volume factor (X_c) as shown in Fig. 3c was derived from the data using the relation $X_c = I_{518} / (I_{518} + I_{480})$, in which I_{518} and I_{480} are the integrated areas extracted from Gaussian fitting at 518 and 480 cm^{-1} , respectively. We can see in Fig. 3a that the n- $\mu\text{-SiO}_x$ layer structure without a bottom-seed-layer appears to present a clear crystalline phase, which is characterised by a sharp Raman peak at 518 cm^{-1} , mainly visible at thicknesses above the value of 17 nm. On the contrary, for thicknesses in the range of 5–13 nm, the structure can be regarded essentially as an amorphous phase, as defined by a broad spectral range centered around 480 cm^{-1} . It has been known that the growth of crystalline phase in silicon films during plasma deposition is generally initiated through a first incubation amorphous phase followed by nucleation and, finally, crystalline phase formation [27–29]. In this process, nucleation-growth aggregation plays an important role to prompt subsequent crystalline particles. As shown in Fig. 3b, the crystalline phase of n- $\mu\text{-SiO}_x$ layer can be initiated even at a low thickness of 7 nm in conjunction with a previously

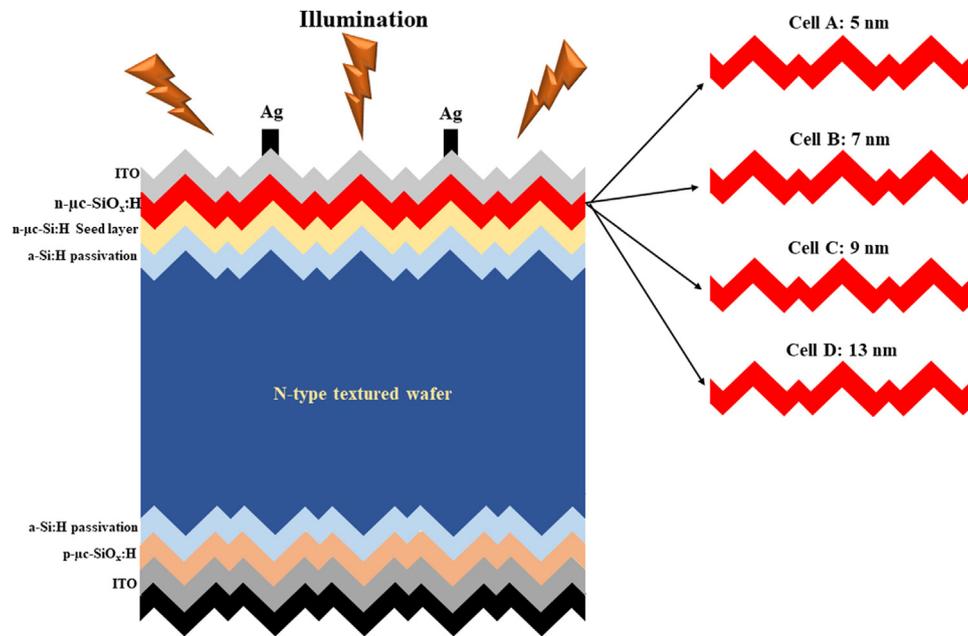


Fig. 1. Schematic diagram RE-SHJ solar cell structure with various thickness of $n\text{-}\mu\text{c-SiO}_x\text{:H}$ front contact layers.

prepared seed layer. Table 1 lists the deposition parameters of the seed layer prepared in very high $[\text{H}_2]/[\text{SiH}_4]$ gas ratio of 160. It has been reported by many research groups that the seed layer formed at very high $[\text{H}_2]/[\text{SiH}_4]$ gas ratios (~ 160 , without added CO_2 gas content) has a hydrogenated silicon structure with micro or nano-crystallite particles including vertical columnar-crystalline phases enveloped by amorphous phase at the boundaries [27,30–32]. Here, because of the seed layer thickness of approximately 3 nm, a crystalline-nucleation behaviour is observed for subsequent layer growth. The seed layer plays a significant role in prompting the rapid growth of crystalline phase in the $n\text{-}\mu\text{c-SiO}_x$ layer even with a low thickness of 7 nm. Based on the nucleation kinetics, the crystalline phase may grow rapidly and increase in thickness [28]. This is discernible in Fig. 3c where the value of X_c increases with thickness of $n\text{-}\mu\text{c-SiO}_x\text{:H}$ layer.

The crystalline phase formation is considerably meaningful when high conductivity of bi-phase front $n\text{-}\mu\text{c-SiO}_x$ material is regarded. As seen in Table 2, the conductivity of the 7-nm thick $n\text{-}\mu\text{c-SiO}_x$ -seed layer stack is four orders of magnitude higher than that of the film without the seed layer. The conductivity is seen to increase with thickness. From the behaviour of X_c in Fig. 3c, it infers that the conductivity appears to rise with increase in the crystalline volume factor (X_c). In addition, the sooner crystalline phase formation of films at 7–9 nm thickness with seed layer is also significant meaningful to reducing the thickness of $n\text{-}\mu\text{c-SiO}_x$ front layer, i.e. high transparency and/or low parasitic absorption. We can see that while the conductivity is nearly analogous with 17 nm film without seed layer, the refractive index at 400 nm and 632 nm wavelength of 7–9 nm-film with seed layer is significantly

lower. This means that by maintaining the high conductivity for effective carrier collection to obtain high V_{oc} and FF, the thin $n\text{-}\mu\text{c-SiO}_x$ front layer (thickness of 7–9 nm) together with the seed layer can significantly lower parasitic absorption losses; this yields higher values of J_{sc} than that in the case of using thicker (greater than 17 nm) $n\text{-}\mu\text{c-SiO}_x$ front layers.

To validate the observed properties, several RE-SHJ solar cells with different $n\text{-}\mu\text{c-SiO}_x$ front layer thicknesses were fabricated on 3-nm thick seed layers in a configuration shown in Fig. 1. A RE-SHJ reference cell (cell E) with 17-nm thick $n\text{-}\mu\text{c-SiO}_x$ front layer but without the seed layer was also prepared for control data. The cell parameters, namely V_{oc} , J_{sc} , FF, and efficiency (E_{ff}) of all the cells fabricated are summarized in Table 3 and J-V curves under one-sun illumination were plotted in Fig. 4a. We can see that the lowest V_{oc} and FF values are attributed to a cell (A) which is probably due to incomplete crystallisation of thin $n\text{-}\mu\text{c-SiO}_x$ front layer of 5 nm as indicated in Fig. 3b. With completely formed crystalline phase on the seed layer of the $n\text{-}\mu\text{c-SiO}_x$ front layer, the V_{oc} and FF of the cells (from B to D) are considerably higher than those of cell A. It is notable that with the seed layer on cells B, C, and D, the parameters V_{oc} and J_{sc} have significantly improved in comparison with those of the reference cell (E). The higher values of J_{sc} are attributed to the drastically reduced thickness of the front layer, which can result in significantly lower losses due to parasitic absorption and internal recombination. The former is discernible in the EQE characteristics as shown in Fig. 4b. The EQE of the cell with the seed layer is considerably higher than that of the reference cell, especially in short wavelength region of 300–600 nm.

Table 2

Conductivity and refractive index of $n\text{-}\mu\text{c-SiO}_x\text{:H}$ with various thickness deposited with and without seed layer.

Thickness (nm)	With seed layer		Without seed layer			
	Conductivity (S/cm)	Refractive index at		Conductivity (S/cm)	Refractive index at	
		400 nm	632 nm		400 nm	632 nm
5	2.0×10^{-6}	3.42	2.76	1.0×10^{-6}	2.70	2.34
7	1.5×10^{-2}	3.45	2.79	3.0×10^{-6}	2.83	2.40
9	7.2×10^{-2}	3.48	2.81	9.0×10^{-6}	2.95	2.48
13	1.2×10^{-1}	3.70	2.91	2.1×10^{-4}	3.30	2.7
17	9.0×10^{-1}	3.72	3.0	8.0×10^{-2}	3.55	2.85

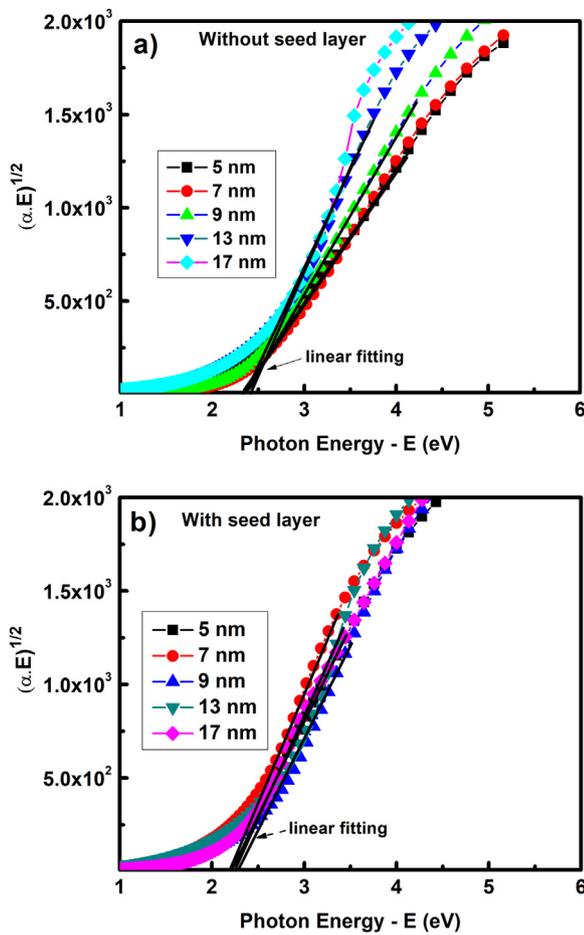


Fig. 2. Plots of $\{\alpha(E)E\}^{1/2}$ versus photon energy E from which the optical band gap is extracted by the linear fitting.

The observed high V_{oc} of the cells B, C and D can be understood from the contribution of the $n-\mu\text{C-Si:H}$ seed layer. Mazzarella et al. reported that the surface passivation quality of the wafer improves during growth of subsequent $n-\mu\text{C-Si:H}$ layer in very high hydrogen dilution [33]. This is akin to the post-hydrogen plasma treatment which is often applied to a very thin $a\text{-Si:H}$ layer to passivate dangling bonds on the wafer surface. To understand the mechanism better, we measured the minority carrier lifetime in all cell precursor components prior to ITO deposition, i.e. after p- and n-layer deposition as shown in Fig. 5a. Measured values were compared with that of wafer prior to n-layer deposition (reference wafer - RW) as shown in Fig. 5b. In addition, Fig. 5c summarizes the lifetime and implied V_{oc} ($i-V_{oc}$) values at one-sun illumination of the components as depicted in Fig. 5a. Fig. 5c shows that the lifetime and $i-V_{oc}$ of the reference cell (E) were insignificantly higher compared with that of RW while that of the cells (B – D) were considerably enhanced compared with that of RW. This indicates that the seed $\mu\text{C-Si:H}$ layer played essentially role in not only promoting fast the crystalline-growth of subsequent $\mu\text{C-SiO}_x\text{:H}$ layer but also enhancing passivation quality of wafer surface. As a results, the V_{oc} of cells B-D was considerably higher than that of cell E, as shown in Table 3. A lower lifetime and $i-V_{oc}$ of cell A compared with that of cell E is probable to the uncompleted crystallisation of thin $n-\mu\text{C-SiO}_x$ front layer which can cause more defects in film structure and thus deteriorate the V_{oc} and FF, as shown in Table 3. It can be concluded that if the crystalline-growth is enhanced by the seed layer, the subsequent $n-\mu\text{C-SiO}_x\text{:H}$ layer can gain a sufficient conductivity even at a thin thickness to reduce the parasitic absorption losses. In addition, complete crystalline-growth of $n-\mu\text{C-SiO}_x$ front layer on the seed layer can enhance the passivation quality of wafer surface and thus result in high V_{oc} and

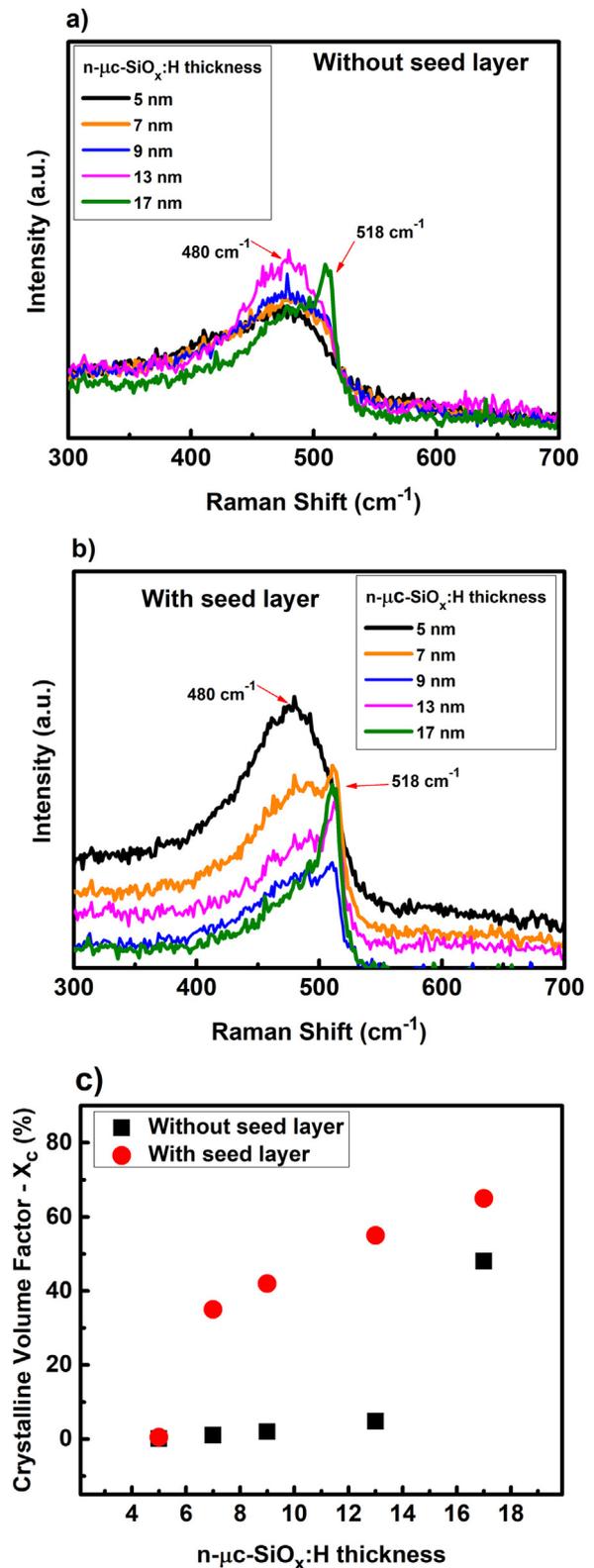


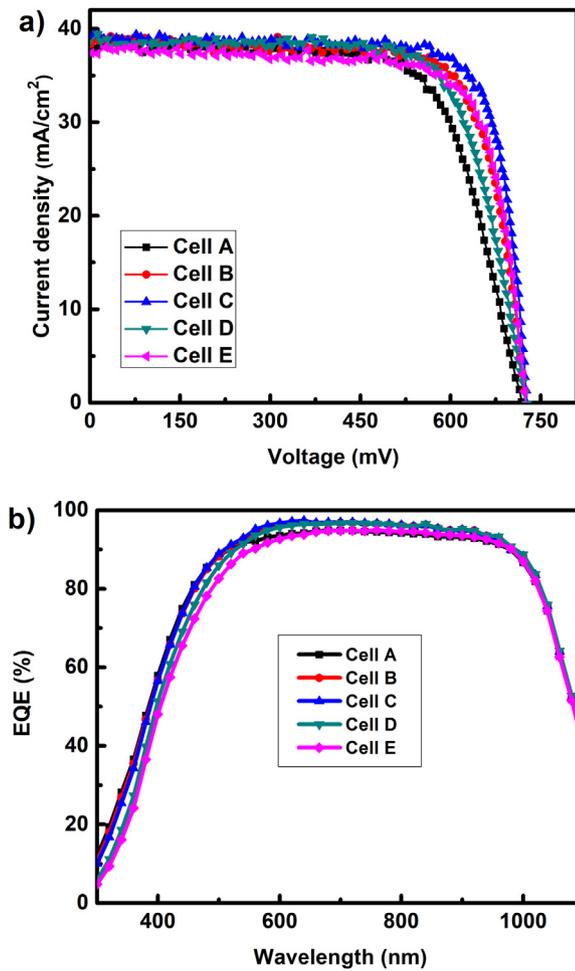
Fig. 3. Raman spectra of $n-\mu\text{C-SiO}_x\text{:H}$ front contact layers: a) without and b) with the seed layer; and c) the crystalline volume factors (X_c) of $n-\mu\text{C-SiO}_x\text{:H}$ front contact layers with and without the seed layer as a function of thickness.

FF.

Basically, a solar cell is attributed to be a diode. In the dark, characteristic of a solar cell behaves completely like a diode. We measured and analyzed the dark I–V (DIV) characteristic on each of the cells based on the diffusion model of the single diode, to obtain better insight

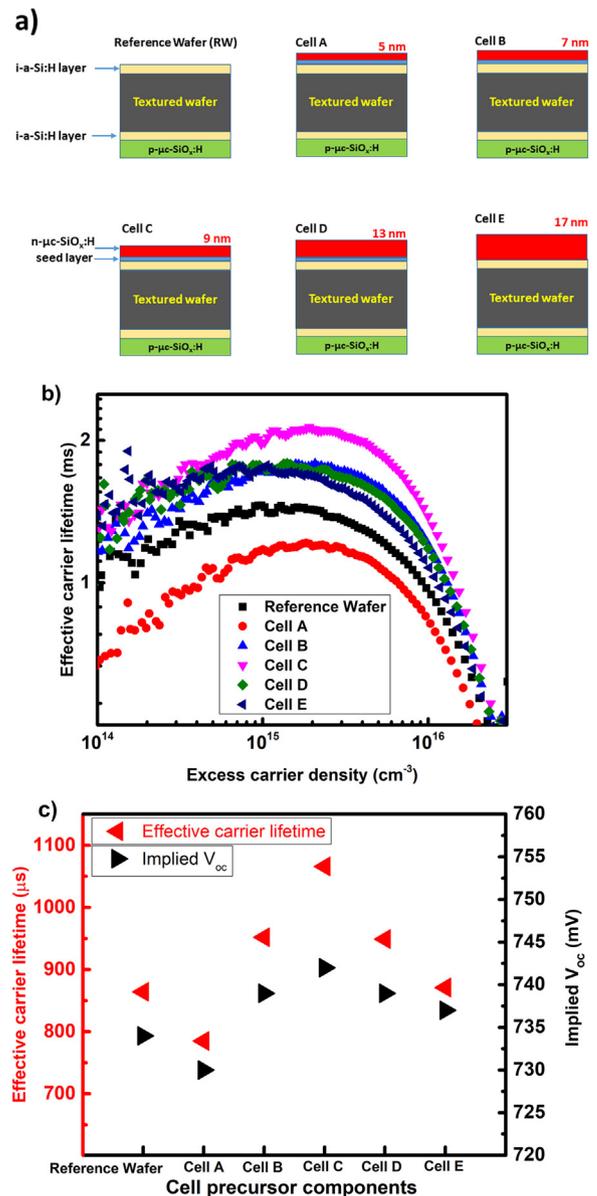
Table 3The cell parameters including V_{oc} , J_{sc} , FF and Efficiency of SHJ solar cells.

RE-SHJ solar cells	Seed layer thickness (nm)	n- μ c-SiO _x :H thickness (nm)	V_{oc} (mV)	J_{sc} (mA/cm ²)	FF (%)	Efficiency (%)
A	3	5	717	38.3	70	19.2
B	3	7	725	39	75	21.2
C	3	9	727	39	77	21.8
D	3	13	726	38.5	75	21
E	0	17	723	38	76	20.8

**Fig. 4.** a) J-V curves under one-sun illumination and b) external quantum efficiencies of RE-SHJ solar cells.

into the enhanced quality of the cells. From DIV characteristics, the diode ideal factor (n_0) and inverse saturation current density (J_0) were extracted as shown in Fig. 6a and b [34] and summarized in Fig. 6c. For crystalline silicon solar cells, n_0 towards 1 means that the diffusion current can be preponderant. In case of n_0 towards a value far higher than 1, this means that the generation/recombination current dominates. Cell A gives the highest values of n_0 and J_0 as shown in Fig. 6c. A higher value of n reflects when the recombination current dominates the diode conduction mechanism [35,36]. It is now clearly evident that cell A is beset with defects caused by incomplete crystallisation. With improved crystallisation and thinner front layers of n- μ c-SiO_x, cells B and D display much lower n_0 and J_0 values than that of the reference cell (E) thus reinforcing the critical role of the seed layer in reducing internal recombination losses in the n- μ c-SiO_x front layer and to realise high values of V_{oc} .

In conclusion, doped wide-gap μ c-SiO_x:H materials have been used successfully as front contact layer in SHJ solar cells [25,33,37]. The

**Fig. 5.** a) Schematic diagram of all cell precursor components for the minority carrier lifetime measurement; b) effective carrier lifetime of these precursor components as a function of excess carrier density; and c) effective carrier lifetime and implied V_{oc} values at one-sun illumination as a function of cell precursor components.

challenge is still to grow very thin and highly crystalline films to reduce significantly the parasitic absorption while maintaining the effective carrier collection. Here, we indicated that the n- μ c-SiO_x:H front layer grown on the n- μ c-Si:H seed layer can gain highly crystalline with considerably thinner thickness than the n- μ c-SiO_x:H without seed layer. It has been indicated that introducing oxygen into μ c-Si:H network restricts the crystalline-growth because numerous Si will bond to

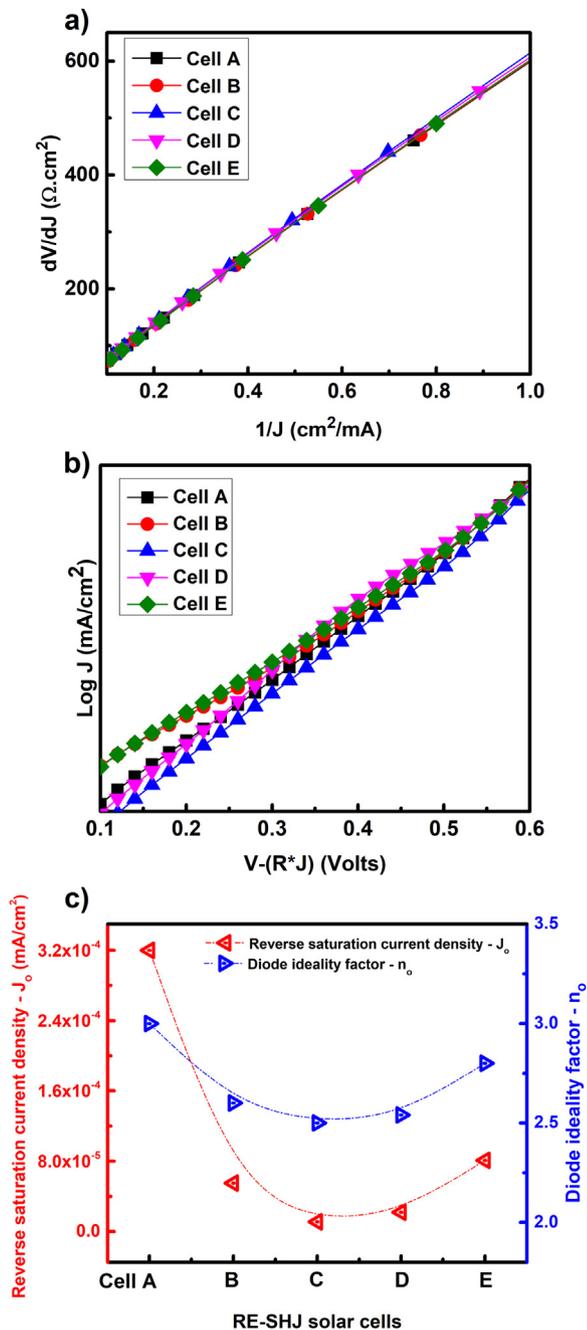


Fig. 6. a) Plot of dV/dJ as a function of $1/J$ for all RE-SHJ solar cells (the intercepts of linear fittings are series resistances (R_s)); b) Plot of $\log J$ as a function of $(V - JR)$ (the slopes of linear fittings give q/n_0kT ratios and intercepts are J_0 values); and c) n_0 and J_0 values of all RE-SHJ solar cells.

oxygen into Si-O-Si form, and thus diminishing the possible number of available Si to form Si-Si bonding in nano-crystallites [26,37,38]. In addition, $\mu\text{-Si:H}$ film quality deteriorates with incorporating oxygen [39]. Consequently, it will be better if using a $\mu\text{-Si:H}$ materials as an initial seed layer for subsequent crystalline-growth of $n\text{-}\mu\text{-SiO}_x\text{:H}$ layer.

4. Conclusion

We have investigated the optoelectronic properties and structure of $n\text{-}\mu\text{-SiO}_x$ front contact layers of various thicknesses grown on $n\text{-}\mu\text{-Si:H}$ seed layer for RE-SHJ solar cells. At a threshold thickness, the stack (film with a seed layer) shows excellent optoelectronic properties with low refractive index (i.e. high transparency) while maintaining a high

fraction of crystalline phase and high electrical conductivity in comparison to films without the seed layer. In application to RE-SHJ solar cells, the lower refractive index yields J_{sc} improvement by 2.64% which is attributed to reduced parasitic absorption loss by virtue of the low thickness of the front contact layer. The enhancement of 0.56% in V_{oc} is due to improved surface passivation under specific conditions of PECVD growth of the seed layer in very high hydrogen dilution which play a critical role during post-hydrogen plasma treatment. An extremely low thickness of the $n\text{-}\mu\text{-SiO}_x$ front contact layer facilitates reduced internal recombination loss. With optimised seed and front contact layer thickness, the best fabricated RE-SHJ solar cell shows a high efficiency value of 21.8%.

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